

1. Kesiti Klein-Gordonova jednačina (za slobodnu česticu)

$$\left(\square + \frac{m_0^2 c^2}{\hbar^2} \right) \phi(x) = 0 \quad [m_0 = m] \quad (1)$$

Particularno rešenje (1) tražimo u obliku ravnog talasa

$$e^{-ik \cdot x} = e^{-\frac{i}{\hbar} p_\mu x^\mu} = e^{-\frac{i}{\hbar} (p_0 x^0 - \vec{p} \cdot \vec{x})}$$

$$= e^{\frac{i}{\hbar} (\vec{p} \cdot \vec{x} - E t)}$$

(2)

Provera

$$\square = \partial_\mu \partial^\mu = \frac{\partial^2}{\partial x^0{}^2} - \nabla^2$$

$$\left(\frac{\partial^2}{\partial x^0{}^2} - \nabla^2 + \frac{m^2 c^2}{\hbar^2} \right) e^{\frac{i}{\hbar} (\vec{p} \cdot \vec{x} - p_0 x^0)} = 0$$

$$e^{\frac{i}{\hbar} (\vec{p} \cdot \vec{x} - p_0 x^0)} \left(\frac{i}{\hbar} p_0 \right)^2 - e^{\frac{i}{\hbar} (\vec{p} \cdot \vec{x} - p_0 x^0)}$$

$$\left\{ \left(\frac{i}{\hbar} p_1 \right)^2 + \left(\frac{i}{\hbar} p_2 \right)^2 + \left(\frac{i}{\hbar} p_3 \right)^2 \right\} + \frac{m^2 c^2}{\hbar^2} e^{\frac{i}{\hbar} (\vec{p} \cdot \vec{x} - p_0 x^0)} = 0$$

$$\left(-\frac{p_0^2}{\hbar^2} + \frac{p_1^2}{\hbar^2} + \frac{p_2^2}{\hbar^2} + \frac{p_3^2}{\hbar^2} + \frac{m^2 c^2}{\hbar^2} \right) e^{\frac{i}{\hbar} (\vec{p} \cdot \vec{x} - p_0 x^0)} = 0$$

$$p_0 = \frac{E}{c}$$

$$-p_0^2 + \vec{p}^2 + m^2 c^2 = 0$$

$$-\frac{E^2}{c^2} + \vec{p}^2 + m^2 c^2 = 0 \quad / \cdot c^2$$

$$-E^2 = -\vec{p}^2 c^2 - m^2 c^4, \quad \text{ty}$$

$$\boxed{E^2 = c^2 \vec{p}^2 + m^2 c^4}$$

Dakle, $e^{-i\vec{k}\cdot\vec{x}}$ jeste resenje KG j-ne.

Ali

$$E = \pm \sqrt{c^2 \vec{p}^2 + m^2 c^4}$$

Postoje i negativna resenja!

Za fiksnu vrednost vektora \vec{k} ,
postoje dva linearno nezavisna resenja
 $e^{-i\omega_{\vec{k}}t + i\vec{k}\cdot\vec{x}}$ i $e^{i\omega_{\vec{k}}t + i\vec{k}\cdot\vec{x}}$
B. 203200

Oste resenje j-ne (I) je

$$\phi(x) = \frac{1}{(2\pi)^{3/2}} \int \frac{d^3\vec{k}}{\sqrt{2\omega_{\vec{k}}}} \left(a(\vec{k}) e^{-i(\omega_{\vec{k}}t - \vec{k}\cdot\vec{x})} + b^\dagger(-\vec{k}) e^{i(\omega_{\vec{k}}t + \vec{k}\cdot\vec{x})} \right)$$

$a(\vec{k})$
 $b^\dagger(\vec{k})$ } kompleksni koeficijenti

Pogledati Mušicki

III/2 (44.14)



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Uz smenu

$$\vec{k} \rightarrow -\vec{k}$$

u drugom članu

$$\phi(x) = \frac{1}{(2\pi)^{\frac{3}{2}}} \int \frac{d^3\vec{k}}{\sqrt{2\omega_k}} \left(a(\vec{k}) e^{-i\vec{k}\cdot\vec{x}} + b^\dagger(\vec{k}) e^{i\vec{k}\cdot\vec{x}} \right)$$

postizemo se komplementarni zapis!

Решение о котором не речу оговорку
 свойствами функции комплексной, т.е. чисте
 св. базис (т.е. оговорку полярной и негравитной
 энергии)

$$\int e^{-ikx} e^{ikx'} d^4k = \int e^{ik(x'-x)} d^4k$$

$$= (2\pi)^4 \delta^4(x'-x)$$

Како не у пазару ОДБ \Rightarrow у пазару не ЛНЗ скан

ЛНЗ система ф-та сь може утврдити и
 преко Вронскиана

$$\begin{vmatrix} f_1(x) & f_2(x) \\ f_1'(x) & f_2'(x) \end{vmatrix} = \begin{vmatrix} e^{ikx} & e^{-ikx} \\ ik e^{ikx} & -ik e^{-ikx} \end{vmatrix} = 2ik \neq 0$$

$k \neq 0$

$$W = 0 \not\Rightarrow \Lambda 3$$

$$\Lambda 3 \Rightarrow W = 0$$

Đata je struja

$$\dot{j}_\mu = i (\phi \partial_\mu \phi^* - \phi^* \partial_\mu \phi)$$

Pokazati da ona zadovoljava i-nu kontinuiteta

$\partial^\mu \dot{j}_\mu = 0$, $\phi = \phi(x)$ se bje K G i-n E .

$$\partial^\mu \dot{j}_\mu = i (\partial^\mu \phi \partial_\mu \phi^* + \phi \partial^\mu \partial_\mu \phi^* - \partial^\mu \phi^* \partial_\mu \phi - \phi^* \partial^\mu \partial_\mu \phi)$$

$$= i (\underbrace{\partial^\mu \phi \partial_\mu \phi^*}_{\square \phi \phi^*} + \phi \square \phi^* - \underbrace{\partial^\mu \phi^* \partial_\mu \phi}_{\square \phi^* \phi} - \phi^* \square \phi)$$

$$(\square + m^2) \phi(x) = 0$$

$$(\square + m^2) \phi^*(x) = 0$$

$$\partial^\mu \dot{j}_\mu = 0$$

$$\dot{j}_\mu = (c\mathcal{E}, -\vec{j})$$

Ali,

$$\dot{j}_0 = i (\phi \partial_0 \phi^* - \phi^* \partial_0 \phi) = c\mathcal{E}$$

Onda \mathcal{E} može biti i negativno,

nasuprot nerelativističnoj kvantnoj mehanici

gde je $\mathcal{E} = |\Psi|^2 \geq 0$.

Razlog: K G jednačina je drugog reda po vremenu!

3. Koristeci kanonske komutacione relacije

$$[\phi(\vec{x}, t), \dot{\phi}(\vec{y}, t)] = i\delta^{(3)}(\vec{x} - \vec{y})$$

$$[\phi(\vec{x}, t), \phi(\vec{y}, t)] = [\dot{\phi}(\vec{x}, t), \dot{\phi}(\vec{y}, t)] = 0$$

izvesti sledece komutacione relacije za kreacione i annihilacione operatore

$$[a(\vec{k}), a^\dagger(\vec{l})] = \delta^3(\vec{k} - \vec{l})$$

$$[a(\vec{k}), a(\vec{l})] = [a^\dagger(\vec{k}), a^\dagger(\vec{l})] = 0$$

Razmatraci sledeci realnog KG polja.

Overste referenc KG - j-nB

$$\phi(x) = \frac{1}{(2\pi)^{\frac{3}{2}}} \int \frac{d^3\vec{k}}{\sqrt{2\omega_k}} (a(\vec{k})e^{-ik\cdot x} + b^\dagger(\vec{k})e^{ik\cdot x})$$

Ali se od stvarnih polja razlikuje
~~u tome sto su stvarna polja realna~~

Za realno polje $a(\vec{k}) = b(\vec{k})$

$$\phi = \frac{1}{(2\pi)^{\frac{3}{2}}} \int \frac{d^3\vec{k}}{\sqrt{2\omega_k}} [a(\vec{k})e^{-ik\cdot x} + a^\dagger(\vec{k})e^{ik\cdot x}]$$

$\phi(x)$

$$\pi = \dot{\phi}$$

$$\dot{\phi} = \frac{1}{(2\pi)^{\frac{3}{2}}} i \int \frac{d^3\vec{k}}{\sqrt{2\omega_k}} \omega_k [-a(\vec{k})e^{ik\cdot x} + a^\dagger(\vec{k})e^{ik\cdot x}]$$

$$\int d^3\vec{r} \phi(x) e^{-i\vec{k}'\cdot\vec{r}} =$$

$$\int d^3\vec{r} \left[\int \frac{d^3\vec{k}}{(2\pi)^{3/2} \sqrt{2\omega_k}} \left[a(\vec{k}) e^{-i\vec{k}\cdot\vec{r}} + a^\dagger(\vec{k}) e^{i\vec{k}\cdot\vec{r}} \right] \right]$$

$$e^{i\vec{k}'\cdot\vec{r}} =$$

$$\int d^3\vec{r} \left[\int \frac{d^3\vec{k}}{(2\pi)^{3/2} \sqrt{2\omega_k}} \left[a(\vec{k}) e^{-i(\omega_k t - \vec{k}\cdot\vec{r})} + a^\dagger(\vec{k}) e^{i(\omega_k t - \vec{k}\cdot\vec{r})} \right] \right]$$

$$e^{-i\vec{k}'\cdot\vec{r}} =$$

$$\int d^3\vec{r} \int d^3\vec{k} a(\vec{k}) e^{-i\omega_k t} e^{i\vec{r}\cdot(\vec{k}-\vec{k}')} +$$

$$\int d^3\vec{r} \int d^3\vec{k} a^\dagger(\vec{k}) e^{i\omega_k t} e^{-i\vec{r}\cdot(\vec{k}+\vec{k}')} =$$

$$\int d^3\vec{r} \int d^3\vec{k} a^\dagger(\vec{k}) e^{i\omega_k t} e^{-i\vec{r}\cdot(\vec{k}+\vec{k}')} =$$

Definieren δ fze:

$$\delta(x-x') = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{iB(x-x')} dB \quad (1D)$$

$$= \frac{1}{(2\pi)^{3/2}} \frac{(2\pi)^3}{\sqrt{2\omega_k}} \int d^3\vec{k} a(\vec{k}) e^{-i\omega_k t} \delta^{(3)}(\vec{k}-\vec{k}') +$$

$$\frac{1}{(2\pi)^{3/2}} \frac{(2\pi)^3}{\sqrt{2\omega_k}} \int d^3\vec{k} a^\dagger(\vec{k}) e^{i\omega_k t} \delta^{(3)}(\vec{k}-(-\vec{k}'))$$

Darje,

$$\int d^3\vec{z} \phi(x) e^{-i\vec{k}' \cdot \vec{z}} =$$

$$\frac{(2\pi)^{\frac{3}{2}}}{\sqrt{2\omega_{k'}}} [a(\vec{k}') e^{-i\omega_{k'}t} + a^\dagger(-\vec{k}') e^{i\omega_{k'}t}] \quad (1)$$

Slično, dobija se i za $\dot{\phi}(x)$

$$\int d^3\vec{z} \dot{\phi}(x) e^{-i\vec{k}' \cdot \vec{z}} = i(2\pi)^{\frac{3}{2}} \sqrt{\frac{\omega_{k'}}{2}} [a^\dagger(-\vec{k}') e^{i\omega_{k'}t} + a(\vec{k}') e^{-i\omega_{k'}t}] \quad (2)$$

Iz (1) i (2) se dobija

$$a(\vec{k}) = \frac{1}{(2\pi)^{\frac{3}{2}}} \frac{1}{\sqrt{2\omega_k}} \int d^3\vec{z} e^{i\vec{k}\cdot\vec{x}} [\omega_k \phi(x) + i \dot{\phi}(x)]$$

$$a^\dagger(\vec{k}) = \frac{1}{(2\pi)^{\frac{3}{2}}} \frac{1}{\sqrt{2\omega_k}} \int d^3\vec{z} e^{-i\vec{k}\cdot\vec{x}} [\omega_k \phi(x) - i \dot{\phi}(x)]$$

Sada tražimo:

$$[a(\vec{k}), a^\dagger(\vec{k}')] =$$

$$\frac{i}{2(2\pi)^3} \frac{1}{\sqrt{\omega_k \omega_{k'}}} \int d^3\vec{z} d^3\vec{z}' e^{i(\vec{k}\cdot\vec{x} - \vec{k}'\cdot\vec{z}')} \quad (*)$$

$$\begin{aligned} X &= (ct, \vec{z}) \\ Y &= (ct, \vec{z}') \end{aligned}$$

$$(-\omega_k [\phi(x), \phi(y)] + \omega_{k'} [\phi(x), \phi(y)])$$

$$= \frac{1}{2(2\pi)^3} \frac{1}{\sqrt{\omega_k \omega_{k'}}} \int d^3\vec{r} e^{i(\omega_k - \omega_{k'})t + i(\vec{q}' - \vec{k})\vec{r}}$$

$$(\omega_k + \omega_{k'}) = \underline{\underline{\delta^3(\vec{k} - \vec{q}')}}$$

Za domaći ostatak

$$= \frac{1}{2} \frac{1}{\sqrt{\omega_k \omega_q}} e^{i(\omega_k - \omega_q)t} \int \frac{1}{(2\pi)^3} d^3\vec{r} e^{i(\vec{q} - \vec{k})\vec{r}}$$

$\delta^3(\vec{k} - \vec{q})$

$\omega_k = \omega_q$

$$\left(\frac{\otimes}{\otimes}\right) \frac{1}{2} \frac{1}{\omega_k} e^{0} 2\omega_k = 1 \quad \checkmark$$

4. Pokazati da za vertičnu spina nula i nabeltrisanyu Σ u spohyryem EM pohyry stryca glasi:

pri čemu je ϕ rešeny KG i-ny y spohyryem EM pohyry potenyajala A^μ .

$$j^\mu = (\phi^* \partial^\mu \phi - \phi \partial^\mu \phi^*) + 2iq \phi^* A^\mu \phi$$

Minimalna preskrycaja

$$D_\mu = \partial_\mu + iq A_\mu$$

\Downarrow

$$\bar{D}_\mu = \partial_\mu - iq A_\mu$$

$$D^\mu = \partial^\mu + iq A^\mu$$

$$\bar{D}^\mu = \partial^\mu - iq A^\mu$$

$(D_\mu D^\mu - m^2) \phi = 0$ (1). Pa onda $\phi^*/(1)$ a potom

$$\left. \begin{aligned} \phi^* (D_\mu D^\mu - m^2) \phi &= 0 \\ \phi (\bar{D}_\mu \bar{D}^\mu - m^2) \phi^* &= 0 \end{aligned} \right\} (1)^* \text{ i } \phi/(1)^*$$

$$\left. \begin{aligned} \phi^* D_\mu D^\mu \phi &= m^2 \phi^* \phi \\ \phi \bar{D}_\mu \bar{D}^\mu \phi^* &= m^2 \phi \phi^* \end{aligned} \right\} -$$

$$\phi^* D_\mu D^\mu \phi - \phi \bar{D}_\mu \bar{D}^\mu \phi^* = 0$$

$$\phi^* (\partial_\mu + i g A_\mu) (\partial^\mu + i g A^\mu) \phi - \phi (\partial_\mu - i g A_\mu) (\partial^\mu - i g A^\mu) \phi^* = 0$$

$$\phi^* (\partial_\mu \partial^\mu \phi + i g \partial_\mu (A^\mu \phi) + i g A_\mu \partial^\mu \phi - g^2 A_\mu A^\mu \phi) -$$

$$- \phi (\partial_\mu \partial^\mu \phi^* - i g \partial_\mu (A^\mu \phi^*) - i g A_\mu \partial^\mu \phi^* - g^2 A_\mu A^\mu \phi^*) = 0$$

$$\phi^* \partial_\mu \partial^\mu \phi + i g \phi^* (\partial_\mu A^\mu) \phi + i g \phi^* A^\mu \partial_\mu \phi + i g \phi^* A_\mu \partial^\mu \phi -$$

$$- g^2 \phi^* A_\mu A^\mu \phi - \phi \partial_\mu \partial^\mu \phi^* + i g \phi (\partial_\mu A^\mu) \phi^* + i g \phi A^\mu \partial_\mu \phi^* +$$

$$+ i g \phi A_\mu \partial^\mu \phi^* + g^2 \phi A_\mu A^\mu \phi^* = 0$$

$$(\phi^* \partial_\mu \partial^\mu \phi - \phi \partial_\mu \partial^\mu \phi^*) + 2i g \phi \phi^* (\partial_\mu A^\mu) + 2i g \phi^* A^\mu \partial_\mu \phi +$$

$$+ 2i g \phi A^\mu \partial_\mu \phi^* = 0$$

$$(\phi^* \partial_\mu \partial^\mu \phi - \phi \partial_\mu \partial^\mu \phi^*) + 2i g \phi \phi^* (\partial_\mu A^\mu) +$$

$$+ 2i g A^\mu (\phi^* \partial_\mu \phi + \phi \partial_\mu \phi^*) = 0$$

$$\partial_\mu (\phi \phi^*)$$

$$(\phi^* \partial_\mu \partial^\mu \phi - \phi \partial_\mu \partial^\mu \phi^*) + 2i g \phi \phi^* (\partial_\mu A^\mu) + 2i g A^\mu \partial_\mu (\phi \phi^*) = 0$$

$$(\phi^* \partial_\mu \partial^\mu \phi - \phi \partial_\mu \partial^\mu \phi^*) + 2i g \partial_\mu (\phi \phi^* A^\mu) = 0$$

$$\partial_\mu (\phi^* \partial^\mu \phi - \phi \partial^\mu \phi^*) + 2i g \partial_\mu (\phi \phi^* A^\mu) = 0$$

$$\partial_\mu \left[(\phi^* \partial^\mu \phi - \phi \partial^\mu \phi^*) + 2i g \phi \phi^* A^\mu \right] = 0 \quad \rightarrow \text{j}^\mu \text{ (struja)}$$